

# Sensors

Note: This section is currently under revision.

Sensors produce measurements that reflect the state of the external world - plus some noise. This section will go through what various sensors are measuring, their sources of noise, and how to relate them to states we want to keep track of.

Towards the end, a large example Emission Matrix will be given that demonstrates how to integrate several sensors into a Kalman Filter.

## Accelerometer

An Accelerator measures acceleration. It's traducer is tantamount to a heavy mass in a box, connected to each wall by a spring. As the box gets accelerated from a direction, the mass's inertia resists, so the spring on one side is compressed, and the spring on the other is expanded. By measuring this contraction/expansion, the acceleration is estimated.

Modern microchip accelerometers are thin pieces of silicon, produced with similar fabrication techniques as microprocessors. The motion of the silicon chip is measured by changing the distance in several capacitors, which measurably alters the capacitance. However, the simpler mass-spring conceptualization is easier to talk about, so we'll stick with that.

The accelerometer feels all forces. This include gravity. If you push on a box, the mass will [relatively speaking] move in the opposite direction of acceleration. However, if you place the box on a table, gravity will cause it to sag, compressing the bottom spring and extending the top spring. The accelerometer will interpret it as being pushed *upwards*, not downwards. As such, we can see that *body* forces, like gravity, are not actually felt by the accelerometer. Only normal forces generated in response to them.

All together, an accelerometer feels 4 different forces. Gravity, Centripetal Acceleration, Rotational Acceleration, and Translational Acceleration.

$$A = a_t + a_r + a_c + a_g$$

## Gravity

The accelerometer feels all forces. This include gravity. If you push on a box, the mass will [relatively speaking] move in the opposite direction of acceleration. However, if you place the box on a table, gravity will cause it to sag, compressing the bottom spring and extending the top spring. The accelerometer will interpret it as being pushed *upwards*, not downwards. As such, we can see that *body* forces, like gravity, are not actually felt by the accelerometer. Only normal forces generated in response to them.

We can safely assume that gravity is always pulling *downwards*, by definition. However, it will not

always add a static acceleration vector along the positive z-axis, because our sensor does not stay level. In fact, measuring gravity gives you an indication of an external static vector. The measurement of this vector is affected by your orientation. Conversely, by measuring your relation to this vector, you can help determine your orientation.

It should be noted that, due to the other sources of acceleration, utilizing an accelerometer's measurements to define your z-axis is likely to be very noisy.

So sitting still at orientation  $R(\psi, \phi, \theta) = R(0, 0, 0)$  the accelerometer would measure  $[0, 0, +g]$  for  $g \approx 9.8 \frac{m}{s^2}$ . If pitched upwards at  $45^\circ$ , the gravity vector would be detected by the positive z and positive x axis at  $45^\circ$ . This transformation would be tantamount to a *negative* pitch forward by the gravity acceleration relative to the accelerometer. As such, the *Inverse* Rotation Matrix  $R^{-1} = R'$  describes how the gravity vector is rotated due to the submarine's orientation.

$$a_g = R^T \begin{bmatrix} 0 \\ 0 \\ g \end{bmatrix} = \begin{bmatrix} 0 & 0 & g \end{bmatrix} R$$

## Centripetal Acceleration

The centripetal acceleration felt by the accelerometer is give by the equation  $\frac{v^2}{r} = \omega^2 r = \dot{\theta}^2 r$ . In 3 dimensions, we need to consider how this force is felt from rotation about each axis.

We can consider an accelerometer aligned with the submarine's internal  $x, y, z$  axis located at Position  $[p_x, p_y, p_z] = [1, 1, 0]$ . A rotation of  $\dot{\theta} = 0.5 \frac{rad}{sec}$  about the z-axis (yaw) would constitute a rotational velocity. The radius  $r = \sqrt{x^2 + y^2} = \sqrt{2}$ . So the centripetal force  $a_{c:yaw} = (0.5^2)(\sqrt{2}) = \frac{1}{4}\sqrt{2} = \sqrt{\frac{2}{4}} = 1 \frac{m}{s^2}$ .

We can also find that the z-position  $p_z$  of the submarine would not impact the centripetal acceleration felt by the accelerometer. This calculated value would hold for any  $[1, 1, p_z]$ .

However, how will this force be detected by the accelerometer? Centripetal Acceleration is another body force. It is accelerating all masses towards the outside - the accelerometer then feels the *inward* normal force created by the rigid structure it's attached to preventing it from moving in the face of the acceleration. So the acceleration vector measured will be negative, pointing back at the origin.

Furthermore, the acceleration felt will be broken up across sensors. Some proportion of it will be felt as negative acceleration along the x-axis, and some proportion along the y-axis. Consider instead if the accelerometer was at the position  $[ \sqrt{2}, 0, 0 ]$ . The centripetal acceleration felt would be the same, but it would be measured entirely as a positive x component  $a_{c:yaw} = [-\sqrt{2}, 0, 0]$ .

Were it at position  $[0, \sqrt{2}, 0]$  the entirety of the acceleration would be felt as a y componenet  $a_{c:yaw} = [0, \sqrt{2}, 0]$ .

We can generalize the solution and calculate the centripetal force felt from roll, pitch, and yaw rotational velocity from an arbitrary  $[x, y, z]$  position.

The position does not influence the rotational velocity - all rotations are felt equally regardless of position on rigid bodies. This leaves the radius to be calculated. The radius  $r = \sqrt{p_x^2 +$

$p_y^2$ }. The angle of this force acts on the accelerometer is also the angle formed from the center of rotation (origin)  $\theta = \tan\{\frac{p_y}{p_x}\}$ .

The math works out quite conveniently.

$$r = \sqrt{p_x^2 + p_y^2}, \quad \theta = \tan\{\frac{p_y}{p_x}\} \quad p_x = r\cos\{\theta\}, \quad p_y = r\sin\{\theta\} \quad a_x = \omega^2 r\cos\{\theta\}, \quad a_y = \omega^2 r\sin\{\theta\}$$

Therefore:

$$a_x = \omega^2 p_y, \quad a_y = \omega^2 p_x, \quad a_z = 0$$

The trigonometric terms cancel out and the acceleration detected along each axis is proportional to the position along that axis. Therefore, we've calculated the  $a_c$  contribution from yaw. We can generalize this calculation to roll and pitch.

## Yaw

$$r = \sqrt{p_x^2 + p_y^2}, \quad \theta = \tan\{\frac{p_y}{p_x}\} \quad p_x = r\cos\{\theta\}, \quad p_y = r\sin\{\theta\} \quad a_x = -\omega^2 r\cos\{\theta\}, \quad a_y = -\omega^2 r\sin\{\theta\} \quad a_x = -p_x * \omega^2, \quad a_y = -p_y * \omega^2, \quad a_z = 0$$

## Pitch

$$r = \sqrt{p_x^2 + p_z^2}, \quad \theta = \tan\{\frac{p_z}{p_x}\} \quad p_x = r\cos\{\theta\}, \quad p_z = r\sin\{\theta\} \quad a_x = \omega^2 r\cos\{\theta\}, \quad a_z = \omega^2 r\sin\{\theta\} \quad a_x = p_x * \omega^2, \quad a_y = 0, \quad a_z = p_z * \omega^2$$

## Roll

$$r = \sqrt{p_y^2 + p_z^2}, \quad \theta = \tan\{\frac{p_z}{p_y}\} \quad p_y = r\cos\{\theta\}, \quad p_z = r\sin\{\theta\} \quad a_y = \omega^2 r\cos\{\theta\}, \quad a_z = \omega^2 r\sin\{\theta\} \quad a_x = 0, \quad a_y = p_y * \omega^2, \quad a_z = p_z * \omega^2$$

## Total Summation

Putting this all together we get the total centripetal acceleration felt on an accelerometer at position  $[p_x, p_y, p_z]$  with rotational velocities  $\dot{R} = [\dot{\psi}, \dot{\phi}, \dot{\theta}]$ .

$$a_x = (0 * \dot{\psi}^2) + (-p_x * \dot{\phi}^2) + (-p_x * \dot{\theta}^2) \quad a_y = (-p_y * \dot{\psi}^2) + (0 * \dot{\phi}^2) + (-p_y * \dot{\theta}^2) \quad a_z = (-p_z * \dot{\psi}^2) + (-p_z * \dot{\phi}^2) + (0 * \dot{\theta}^2)$$

Thus we can turn this into a matrix format and get:

$$a_c = \begin{bmatrix} a_{cx} \\ a_{cy} \\ a_{cz} \end{bmatrix} = \begin{bmatrix} 0 & -p_x & -p_x \\ -p_x & -p_y & 0 \\ -p_y & -p_z & -p_z \end{bmatrix} \begin{bmatrix} \dot{\psi}^2 \\ \dot{\phi}^2 \\ \dot{\theta}^2 \end{bmatrix}$$

$$\dot{\theta}^2 \end{bmatrix} \quad$$

Finally, it's worth noting that because the accelerations felt are a function of rotational velocity squared, the direction of rotation is irrelevant. The force felt is inwards, and favors neither direction. We know this about centripetal force intuitively, but it's nice to see math reflected our common-sense.

## Translational and Rotational Acceleration

The Translational and Rotational Accelerations are both felt as individual point-accelerations. Rotational Acceleration only makes sense as the difference of accelerations felt on two parts of a rigid body. The translational acceleration of the body is the average of these accelerations.

Assume two accelerometers are placed at locations  $p_1[1,0,0]$  and  $p_2[-1,0,0]$ . An acceleration of  $1\frac{m}{s^2}$  along the x-axis would be measured as  $a_1[1,0,0]$  and  $a_2[1,0,0]$ . They would be felt equally. A yaw acceleration  $\alpha_{yaw} = \frac{a}{r} = 1\frac{rad}{sec^2}$  about the origin would be felt as a sideways acceleration on each accelerometer. The forward accelerometer would be accelerated to the left, into the positive y-axis. The back accelerometer would be accelerated in the opposite direction, into the negative y-axis. Thus the accelerations felt would be  $a_1[0,1,0]$  and  $a_2[0,-1,0]$ . They would experience opposite accelerations.

$$a_t = \frac{a_1 + a_2}{2}, \quad a_c = \frac{a_1 - a_2}{2}$$

It should also be noted, that this math describes the translational acceleration felt by the center of rotation. That is defined as the origin, or center of mass, for convenience. But it can technically be defined as any arbitrary set of coordinates. Demonstrating that is an exercise left to the student.

However, this only works for two accelerometers placed along a common axis at equal distances. We want generalized equations for arbitrary accelerometer placement.

Consider if the two accelerometers were instead placed at  $p_1[0,1,0]$  and  $p_2[0,-1,0]$  and subjected to the same translational acceleration  $a_t[1,0,0]$ . The translational accelerations felt would be equal to the first case. Translational acceleration is felt universally on the rigid body. Subjected to the same rotational acceleration of  $\alpha_{yaw} = 1\frac{rad}{sec^2}$ , however, the left accelerometer would be accelerated backwards toward the negative x-axis, with  $a_1[-1,0,0]$  and the right accelerometer would feel an opposite acceleration forward  $a_2[1,0,0]$ . Were the acceleration doubled to  $\ddot{\theta} = 2\frac{rad}{sec^2}$ , the accelerations felt would be equal to  $[-2,0,0]$  and  $[2,0,0]$ .

Like in the centripetal acceleration case, the z-axis position does not affect the rotational forces felt by a yaw acceleration. Using trigonometry, we can find the calculations for acceleration components felt from a given yaw acceleration at an arbitrary point. Again, like with the centripetal force, the rotational force components will depend will depend on angle  $\theta = \tan^{-1}\left(\frac{p_y}{p_x}\right)$ . The difference is the acceleration felt will be orthogonal to that vector, rather than lie alongside it.

$$p = [p_x, p_y, p_z], \quad \alpha = \frac{a}{r} \quad r = \sqrt{p_x^2 + p_y^2}, \quad \theta = \tan^{-1}\left(\frac{p_y}{p_x}\right) \quad p_x = r\cos{\theta}, \quad p_y = r\sin{\theta} \quad a_x = -\alpha * r\sin{\theta}, \quad a_y = \alpha * r\cos{\theta} \quad a_z = 0$$

Note that the acceleration vector is felt orthogonal to the position of the accelerometer. Thus, the

rotational acceleration vector felt from a yaw acceleration is equal to the position vector  $[p_x, p_y, p_z]$  crossed with the z-axis  $[0, 0, 1]$ , scaled by the actual yaw acceleration about the z-axis. This is why we get the  $p_x$  term for the y component, but a negative  $p_y$  term for the x component, and a 0 along the z-axis.

We can take this method of finding the acceleration vector and apply it to pitch and roll accelerations, to find the total rotational acceleration felt.

## Yaw

$$\begin{aligned} p &= [p_x, p_y, p_z], \quad \alpha = \frac{a}{r} \quad r = \sqrt{p_x^2 + p_y^2}, \quad \theta = \tan^{-1}\left(\frac{p_y}{p_x}\right) \\ p_x &= r \cos\{\theta\}, \quad p_y = r \sin\{\theta\} \\ a_x &= -\alpha * r \sin\{\theta\}, \quad a_y = -\alpha * r \cos\{\theta\} \\ a_x &= -p_y * \alpha, \quad a_y = p_x * \alpha, \quad a_z = 0 \end{aligned}$$

## Pitch

$$\begin{aligned} p &= [p_x, p_y, p_z], \quad \alpha = \frac{a}{r} \quad r = \sqrt{p_x^2 + p_z^2}, \quad \theta = \tan^{-1}\left(\frac{p_z}{p_x}\right) \\ p_x &= r \cos\{\theta\}, \quad p_z = r \sin\{\theta\} \\ a_x &= \alpha * r \sin\{\theta\}, \quad a_z = -\alpha * r \cos\{\theta\} \\ a_x &= p_z * \alpha, \quad a_z = -p_x * \alpha, \quad a_y = 0 \end{aligned}$$

## Roll

$$\begin{aligned} p &= [p_x, p_y, p_z], \quad \alpha = \frac{a}{r} \quad r = \sqrt{p_y^2 + p_z^2}, \quad \theta = \tan^{-1}\left(\frac{p_z}{p_y}\right) \\ p_y &= r \cos\{\theta\}, \quad p_z = r \sin\{\theta\} \\ a_y &= -\alpha * r \sin\{\theta\}, \quad a_z = \alpha * r \cos\{\theta\} \\ a_y &= -p_z * \alpha, \quad a_z = p_y * \alpha, \quad a_x = 0 \end{aligned}$$

## Total Summation

Putting this all together we get the total acceleration felt on an accelerometer at position  $[p_x, p_y, p_z]$  with rotational acceleration  $\dot{R} = [\ddot{\psi}, \ddot{\phi}, \ddot{\theta}]$ .

$$\begin{aligned} a_x &= (0 * \dot{\psi}^2) + (p_z * \dot{\phi}^2) + (-p_y * \dot{\theta}^2) \\ a_y &= (-p_z * \dot{\psi}^2) + (0 * \dot{\phi}^2) + (p_x * \dot{\theta}^2) \\ a_z &= (p_y * \dot{\psi}^2) + (-p_x * \dot{\phi}^2) + (0 * \dot{\theta}^2) \end{aligned}$$

Whatever value remains after the rotational components have been removed is the true translational acceleration  $a_t$ . Thus we can turn this into a matrix format and get:

$$\begin{aligned} a &= a_t + a_r = \begin{bmatrix} a_x \\ a_y \\ a_z \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & p_z & -p_y \\ 0 & 1 & 0 & -p_z \\ 0 & p_x & 0 & 0 \\ 0 & 0 & 1 & p_y \\ 0 & -p_x & 0 & 0 \end{bmatrix} \begin{bmatrix} \ddot{x} \\ \ddot{y} \\ \ddot{z} \\ \ddot{\psi} \\ \ddot{\phi} \\ \ddot{\theta} \end{bmatrix} \end{aligned}$$

As we can see, there are six state variables, and we only have a vector of 3 measurements from one accelerometer. Thus the math reflects the realistic constraints - we need at least two accelerometers to separate rotational and translational acceleration components.

## Total Matrix

We can combine everything now to express the entirety of what our accelerometer measures as a function of our state variables.

$$A = a_t + a_r + a_c + gR \quad A = \begin{bmatrix} 1 & 0 & 0 & 0 & p_z & -p_y \\ 0 & 1 & 0 & -p_z & 0 & p_x \\ 0 & 0 & 1 & p_y & -p_x & 0 \end{bmatrix} \begin{bmatrix} \ddot{x} \\ \ddot{y} \\ \ddot{z} \\ \ddot{\psi} \\ \ddot{\phi} \\ \ddot{\theta} \end{bmatrix} + \begin{bmatrix} 0 & -p_x & -p_x \\ 0 & -p_y & -p_z \\ 0 & -p_z & -p_z \end{bmatrix} \begin{bmatrix} \dot{\psi}^2 \\ \dot{\phi}^2 \\ \dot{\theta}^2 \end{bmatrix} + \begin{bmatrix} 0 & 0 & g \end{bmatrix} \begin{bmatrix} R_{11} \\ R_{12} \\ R_{13} \end{bmatrix} + \begin{bmatrix} R_{21} & R_{22} & R_{23} \\ R_{31} & R_{32} & R_{33} \end{bmatrix}$$

Condensing all states into one vector, we get the overall matrix:

$$\begin{bmatrix} 1 & 0 & 0 & 0 & p_z & -p_y & 0 & -p_x & -p_x & 0 & 0 & 0 & 0 & 0 & 0 & 0 & g & 0 & 0 \\ 0 & 1 & 0 & -p_z & 0 & p_x & -p_y & 0 & -p_y & 0 & 0 & 0 & 0 & 0 & 0 & 0 & g & 0 & 0 \\ 0 & 0 & 1 & p_y & -p_x & 0 & -p_z & -p_z & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & g & 0 & 0 \end{bmatrix} \begin{bmatrix} \ddot{x} \\ \ddot{y} \\ \ddot{z} \\ \ddot{\psi} \\ \ddot{\phi} \\ \ddot{\theta} \\ \dot{\psi}^2 \\ \dot{\phi}^2 \\ \dot{\theta}^2 \\ R_{11} \\ R_{12} \\ R_{13} \\ R_{21} \\ R_{22} \\ R_{23} \\ R_{31} \\ R_{32} \\ R_{33} \end{bmatrix}$$

This seems like quite an excessively rectangular matrix to be using. However, other sensors will also utilize some of these states. Thus, this form is preferable, as we can easily integrate it with other sensors simply by expanding the rows of the left matrix.

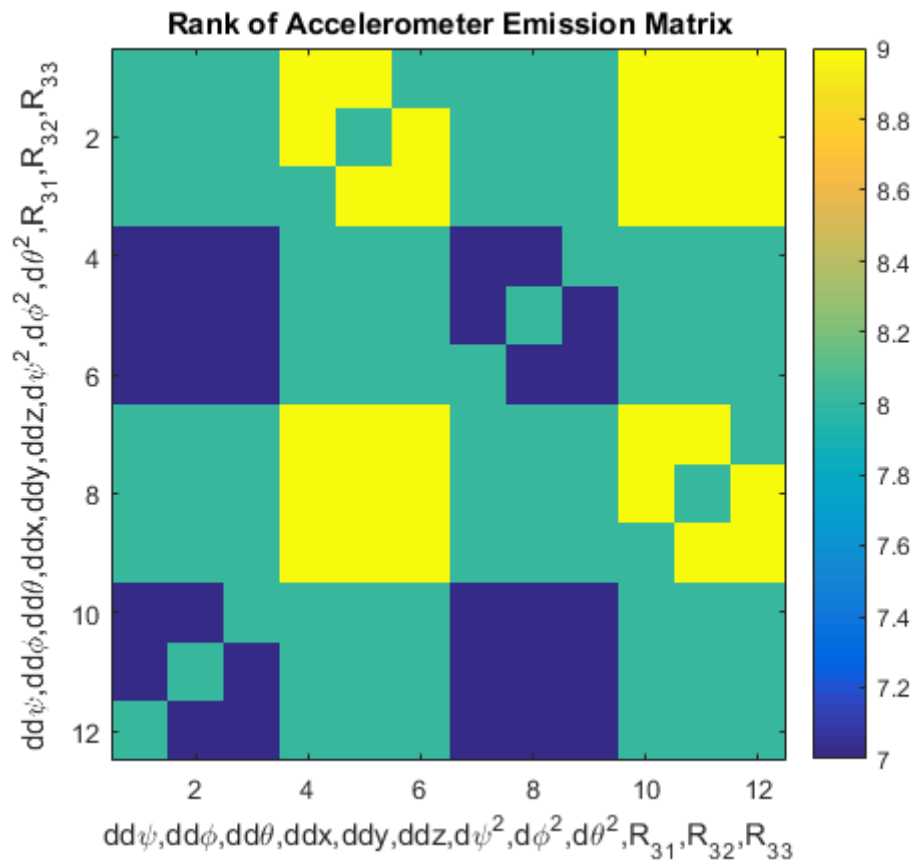
We can also look at our matrix and learn something about what kind of information we need. There are 18 state variables.

6 of them - elements from the  $3 \times 3$  rotation matrix  $R$  - do not interact with our accelerometer at all - which means that accelerometers alone, no matter how many, cannot tell us our orientation. They can tell us, however, the rotation matrix elements  $R_{31}, R_{32}, R_{33}$  which make up the bottom row of the rotation matrix. If you check the 3D Rotation Section in the wiki, you can see that these three pieces of information alone can inform us of our pitch and roll. They can tell us nothing about the yaw of the submarine, which makes perfect sense as detecting a downward pointing gravity vector will never tell us anything about our heading.

Beyond that, there are 9 other values. Naively, we might conclude that with 4 accelerometers placed at four separate locations, we could fully determine the translational acceleration, rotational acceleration, square of the rotational velocity, and the tilt (pitch/roll) of the submarine. All 12 total states.

However, if we use 4 random locations  $p_1, p_2, p_3, p_4$  to construct the matrix, and find its rank, it evaluates to 9. Not 12. It is important the positions be random, so that symmetry does not cheat us out of any information. It may also be worth investigating accelerometers placed precisely at the origin, and unit vectors along each axis, to see if there's any theoretical information available, but realistic constraints prevent us from every placing accelerometers so precisely.

Thus, there are only 9 of those values that can be determined from accelerometers alone. To find which states are linked, we can systematically remove columns from the matrix and re-test its rank. If the rank diminishes, then that column was an independent vector and corresponds to a state not ineffably tied to the value of any other state.



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